

# Nanofluid Flow using Buongiorno Model over a Stretching Sheet and Thermophysical Properties of Nanoliquids

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## Abstract

The nanofluid flow using Buongiorno model over a stretching sheet and thermophysical properties of nanoliquids is numerically studied. The transformation of the governing partial differential equations into nonlinear ordinary differential equations is done using the similarity variables. Therefore, it has been solved numerically using shooting technique. Four nanoparticles (water based fluid) are considered in this study, namely silver Ag, copper Cu, aluminium oxide or known as alumina  $Al_2O_3$  and lastly titanium oxide or known as titania  $TiO_2$ . The numerical results obtained are the velocity profile, temperature profile, and nanoparticle concentration profile. Besides, the results of the local Nusselt number and local Sherwood number are also found. These results are then displayed graphically and argued. The flow and heat transfer performance is discussed in the presence of thermophoresis  $Nt$  and Brownian motion  $Nb$  and nanoparticle volume fraction  $\phi$ . This study has shown that the stretching sheet is an unique solutions. Otherwise, when  $\phi$  increases,  $Nb$  decreases and  $Nt$  decreases, the heat transfer rate increases.

**Keywords:** Boundary Layer, Heat Transfer, Nanofluid, Stagnation Point, Stretching Sheet

## 1. Introduction

Currently, the research about the flow over a stretching plate are investigated by many authors. This is because the applications such as an extrusion, lubricant and glass fiber production are very important. Other applications are wire drawing, metal spinning and hot rolling. The flow past a stretching plate was founded by<sup>1</sup>. Then, this study has been explored by researchers. In<sup>2</sup> figured out the heat transfer over a stretching sheet with permeable surface. Another studies that related to stretching surface are<sup>3-6</sup>. Furthermore, it is noted that<sup>7</sup> was the first to carry out this study in a nanofluid.

The convective heat transfer in nanofluid has been pointed out over the past few years due to the importance of nanofluid. These nanofluid was first introduced by<sup>8</sup>, where the nanoparticles are dispersed in a based fluid. There are many types of nanoparticles that widely been used such as carbon monotubes, metals, carbides or oxides, while the example of based fluids are water, engine oil, ethylene glycol or kerosene. The nanofluid has an unique behaviour that can enhance the thermal conductivity. On the other hand, the nanofluids also been used in many area such as electronics, biomedical, food, transportation and nuclear reactor. Some of the studies that considering nanofluid are<sup>9-11</sup>.

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In<sup>12</sup> introduced nanofluid model which examine the nanofluid flow and heat transfer with the existence of nanoparticles. In<sup>13</sup> used this model to study the hydro-magnetic flow over a permeable sheet. In<sup>14</sup> argued the flow past a stretching or shrinking surface. On the other hand, in<sup>15</sup> introduced the nanofluid model that considers the affects of both Brownian motion parameter and thermophoresis paramater. These two elements are important in nanofluids. A number of researchers have studied this model currently. In<sup>16</sup> studied the flow towards a permeable surface. In<sup>17</sup> investigated the flow over an isothermal vertical plate. Other study, In<sup>18</sup> explained the nanofluid flow over a vertical plate.

We extend and then elaborate the study by<sup>14,16</sup> represent two nanofluids model of<sup>12,15</sup> respectively will be used together as a combination nanofluid equation model in a co-flowing nanofluid. There are very limited study in the literature that deal with these models combination, as mentioned by<sup>19</sup>. We aim to study the presence of thermophoresis, Brownian motion and also the nanoparticle volume fraction. The study also covered the effect of these parameters numerically on both the flow and heat transfer over a stretching plate. As default, the shooting technique will give the results for velocity, temperature and nanoparticle concentration. The results for skin friction coefficient, local Nusselt number and local Sherwood number are also obtained. All these results are shown graphically and elaborated in detail.

## 2. Problem Formulation

The two-dimensional free convection boundary layer flow in the area  $y > 0$ , past a stretching sheet when conditions for the plane,  $y = 0$  and stagnation point,  $x = 0$  are considered. The flow is assumed steady, laminar and incompressible. The flow is assumed steady, laminar and incompressible. Another assumption is the stretching velocity  $U_w(x)$  is vary linearly from  $x = 0$ , with  $U_w(x) = ax$  and  $U_\infty(x) = bx$ , (with  $a$  and  $b$  are both constants where  $b > 0$ ). This assumption also true for the ambient fluid velocity  $U_\infty(x)$ . For stretching sheet,  $a > 0$ , the governing equations are as follows:

$$\frac{\partial u}{\partial x} + \frac{\partial v}{\partial y} = 0, \tag{1}$$

$$u \frac{\partial u}{\partial x} + v \frac{\partial u}{\partial y} = U_\infty \frac{dU_\infty}{dx} + \frac{\mu_{nf}}{\rho_{nf}} \frac{\partial^2 u}{\partial y^2}, \tag{2}$$

$$u \frac{\partial T}{\partial x} + v \frac{\partial T}{\partial y} = \alpha_{nf} \frac{\partial^2 T}{\partial y^2} + \tau \left[ D_B \frac{\partial C}{\partial y} \frac{\partial T}{\partial y} + \left( \frac{D_T}{T_\infty} \right) \left( \frac{\partial T}{\partial y} \right)^2 \right], \tag{3}$$

$$u \frac{\partial C}{\partial x} + v \frac{\partial C}{\partial y} = D_B \frac{\partial^2 C}{\partial y^2} + \left( \frac{D_T}{T_\infty} \right) \frac{\partial^2 T}{\partial y^2}, \tag{4}$$

along the boundary conditions:

$$\begin{aligned} u = U_w(x), \quad v = 0, \quad T = T_w, \quad C = C_w \quad \text{at } y = 0, \\ u \rightarrow U_\infty, \quad T \rightarrow T_\infty, \quad C \rightarrow C_\infty \quad \text{as } y \rightarrow \infty, \end{aligned} \tag{5}$$

with the velocity components  $u$  in  $x$ - direction and the velocity component  $v$  in  $y$ -direction,  $T$ ,  $T_w$  and  $T_\infty$  are temperature, surface temperature, ambient temperature, respectively. Parameters  $C$ ,  $C_\infty$  and  $C_w$  are nanoparticle volume fraction, nanoparticle volume fraction far from the plate and nanoparticle volume fraction at the plate, respectively, thermophoretic diffusion coefficient  $D_T$ , Brownian diffusion coefficient  $D_B$ , heat capacity ratio  $\tau = (\rho C_p)_s / (\rho C_p)_f$ , where nanoparticle heat capacity  $(\rho C_p)_s$  and fluid heat capacity  $(\rho C_p)_f$ . Further,  $\alpha_{nf}$ ,  $\mu_{nf}$  and  $\rho_{nf}$  are the nanofluid thermal diffusivity, nanofluid viscosity and nanofluid density, respectively. These parameters are defined by<sup>20</sup>.

$$\begin{aligned} \alpha_{nf} = \frac{k_{nf}}{(\rho C_p)_{nf}}, \quad \rho_{nf} = (1-\phi)\rho_f + \phi\rho_s, \quad \mu_{nf} = \frac{\mu_f}{(1-\phi)^{2.5}}, \\ (\rho C_p)_{nf} = (1-\phi)(\rho C_p)_f + \phi(\rho C_p)_s, \quad \frac{k_{nf}}{k_f} = \frac{(k_s + 2k_f) - 2\phi(k_f - k_s)}{(k_s + 2k_f) + \phi(k_f - k_s)}. \end{aligned} \tag{6}$$

with parameter nanoparticle volume fraction  $\varphi$ , nanofluid heat capacity  $(\rho C_p)_{nf}$ , nanofluid thermal conductivity  $k_{nf}$ , fluid thermal conductivity  $k_f$ , solids thermal conductivity  $k_s$ , fluid density  $\rho_f$ , solids density  $\rho_s$  and fluid viscosity  $\mu_f$ . Following<sup>21</sup>, the term  $k_{nf}$  is referred as spherical nanoparticles and should not be considered for other shapes. Next, subject to Equation (5), we perform the similarity solution of Equation (1) till Equation (4). The similarity transformations is introduced, see<sup>14,16</sup>:

$$\eta = \left( \frac{U_\infty}{v_f x} \right)^{1/2} y, \quad \psi = (v_f x U_\infty)^{1/2} f(\eta),$$

$$\theta(\eta) = \frac{T - T_\infty}{T_w - T_\infty}, \quad \phi(\eta) = \frac{C - C_\infty}{C_w - C_\infty},$$
(7)

with the similarity variable  $\eta$  and the stream function  $\psi$ , where  $u = \partial\psi / \partial y$  and  $v = -\partial\psi / \partial x$  that satisfies Equation (1). Equation (6) is then substituted into Equation (2) and Equation (3), the transformed equations are:

$$\frac{1}{(1-\varphi)^{2.5} (1-\varphi + \varphi \rho_s / \rho_f)} f''' + ff'' - f'^2 + 1 = 0,$$
(8)

$$\frac{1}{Pr (1-\varphi + \varphi (\rho C_p)_s / (\rho C_p)_f)} \theta'' + \frac{1}{2} f \theta' + Nb \phi' \theta' + Nt \theta^2 = 0,$$
(9)

$$\phi'' + \frac{1}{2} Le f \theta' + \frac{Nt}{Nb} \theta'' = 0,$$
(10)

along with boundary conditions (5), and hence reduced to

$$f(0) = 0, \quad f'(0) = \varepsilon, \quad \theta(0) = 1, \quad \phi(0) = 1,$$

$$f'(\eta) \rightarrow 1, \quad \theta(\eta) \rightarrow 0, \quad \phi(\eta) \rightarrow 0 \quad \text{as } \eta \rightarrow \infty,$$
(11)

with Prandtl number  $Pr = \nu_f / \alpha_f$ , Lewis number  $Le = \nu_f / D_B$  and stretching/shrinking parameter  $\varepsilon = a / b$ , where  $\varepsilon > 0$  for a stretching surface. Parameters  $Nb$  and  $Nt$  are written as

The formulas of physical quantities of interest can be written as

$$Nb = \frac{(\rho C)_p D_B (C_w - C_\infty)}{(\rho C)_f \nu_f}, \quad Nt = \frac{(\rho C)_p D_T (T_w - T_\infty)}{(\rho C)_f T_\infty \nu_f}.$$
(12)

$$C_f = \frac{\tau_w}{\rho_f U_\infty^2}, \quad Nu_x = \frac{x q_w}{k_f (T_w - T_\infty)}, \quad Sh_x = \frac{x q_m}{D_B (C_w - C_\infty)},$$
(13)

Where,

surface shear stress,  $\tau_w = \mu_{nf} \left( \frac{\partial u}{\partial y} \right)_{y=0}$

surface heat flux,  $q_w = -k_{nf} \left( \frac{\partial T}{\partial y} \right)_{y=0}$

mass heat flux,  $q_m = -D_B \left( \frac{\partial C}{\partial y} \right)_{y=0}$ .

(14)

Next, Equation (7) is then substituted into Equations (13) and (14), and therefore the parameters are,

$$C_f Re_x^{1/2} = \frac{1}{(1-\varphi)^{2.5}} f''(0),$$

$$Nu_x Re_x^{-1/2} = -\frac{k_{nf}}{k_f} \theta'(0),$$

$$Sh_x Re_x^{-1/2} = -\phi'(0),$$
(15)

where local Reynolds number  $Re_x = Ux / \nu_f$ .

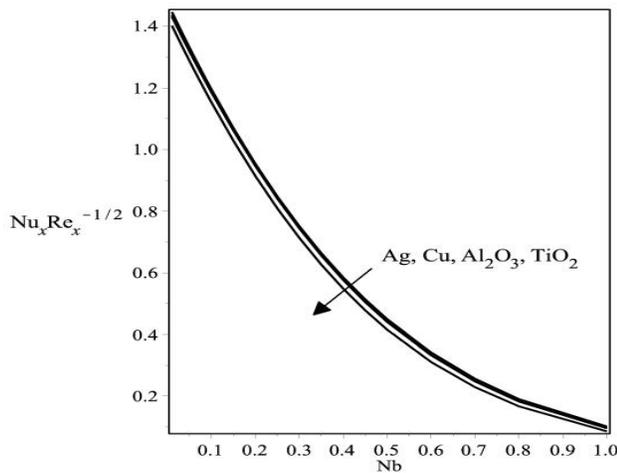
### 3. Results and Discussion

The shooting techniques are used to obtain the numerical solutions for Equation (7) to Equation (9) with the boundary conditions Equation (10). Four different nanoparticles with water based fluid are used to test the influences of  $Nt$ ,  $Nb$ ,  $\varphi$  and  $\epsilon$ . The nanoparticles used are Ag, Cu,  $Al_2O_3$  and  $TiO_2$ . The range of  $\varphi$  should be between 0 to 0.2 where  $\varphi = 0$  is regular fluid<sup>20</sup>. Moreover, the Prandtl number  $Pr$  is 6.2. The properties values of thermophysical are shown in Table 1.

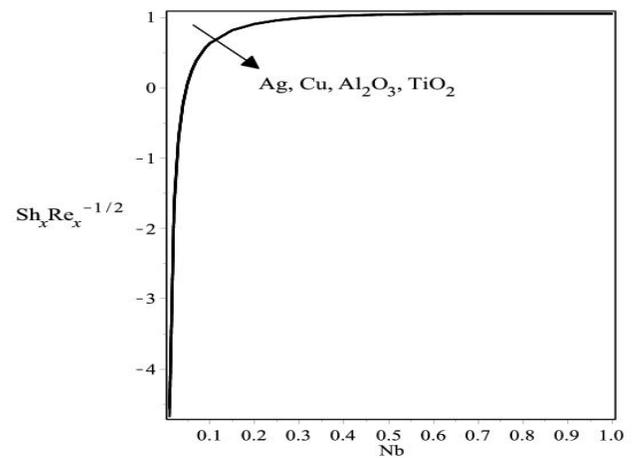
Figures 1 and 2 indicate the variations of local Nusselt number  $Nu_x Re_x^{-1/2}$  local Sherwood number  $Sh_x Re_x^{-1/2}$  for various nanoparticles when  $Pr = 6.2$ ,  $\varphi = 0.1$ ,  $Le = 3$  and  $Nt = 0.1$  with  $\epsilon = 1.2$  for stretching case. It is clearly observed that from Figure 1, the heat transfer rate decreases as  $Nb$  increases while the opposite trend are found in Figure 2. As is observed, the effect of various nanoparticles only gives the small changes for heat and mass flux rates.

**Table 1.** Thermophysical properties of nanofluids<sup>20</sup>

Physical properties	Base fluid				
	Water	Ag	Cu	$TiO_2$	$Al_2O_3$
$C_p (J / kgK)$	4179	235	385	686.2	765
$\rho (kg / m^3)$	997.1	10500	8933	4250	3970
$k (W / mK)$	0.613	429	400	8.9538	40



**Figure 1.** Variation of  $Nu_x Re_x^{-1/2}$  with  $Nb$  for various nanoparticles when  $\epsilon = 1.2$ ,  $Pr = 6.2$ ,  $\varphi = 0.1$ ,  $Le = 3$  and  $Nt = 0.1$ .



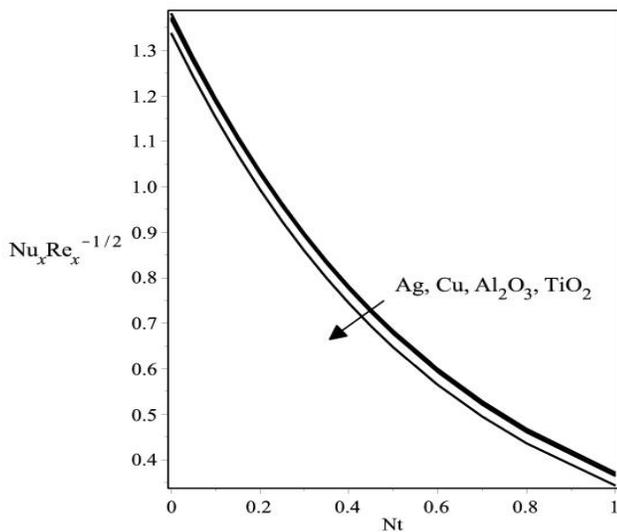
**Figure 2.** Variation of  $Sh_x Re_x^{-1/2}$  with  $Nb$  for various nanoparticles when  $\epsilon = 1.2$ ,  $Pr = 6.2$ ,  $\varphi = 0.1$ ,  $Le = 3$  and  $Nt = 0.1$ .

Moreover, the variation curves of  $Nu_x Re_x^{-1/2}$  and  $Sh_x Re_x^{-1/2}$  with  $Nb$  for various nanoparticles when  $Pr = 6.2$ ,  $\phi = 0.1$ ,  $Le = 3$  and  $Nt = 0.1$  with  $\epsilon = 1.2$  for stretching case are shown in Figure 3 ( $Nu_x Re_x^{-1/2}$ ) and Figure 4 ( $Sh_x Re_x^{-1/2}$ ). From these figures, the local Nusselt number reduces as  $Nb$  increases while the reverse behaviours are found in Figure 4. As is observed, the effect of various nanoparticles only give the small changes for  $Nu_x Re_x^{-1/2}$  and  $Sh_x Re_x^{-1/2}$ .

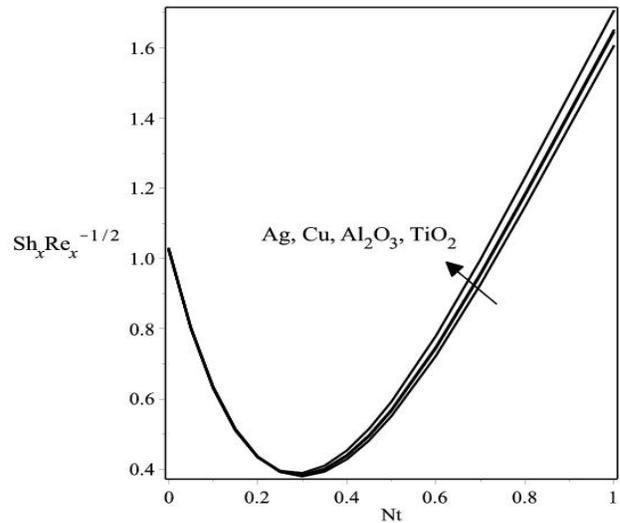
Next, the variations of skin friction coefficient  $C_f Re_x^{1/2}$  together with  $Nu_x Re_x^{-1/2}$  and  $Sh_x Re_x^{-1/2}$  with  $\phi$  for Ag, Cu,  $TiO_2$  and  $Al_2O_3$  when  $\epsilon = 1.2$  (for stretching case) with  $Pr=6.2$ ,  $Nt=0.1$ ,  $Le=3$  and  $Nb=0.1$  are illustrated in Figures 5, 6 and 7, respectively. It is clearly found that, for all these figures show the same trend where  $C_f Re_x^{1/2}$ ,  $Nu_x Re_x^{-1/2}$  and  $Sh_x Re_x^{-1/2}$  increase when  $\phi$  increases. The nanoparticle Ag is the higher value for  $C_f Re_x^{1/2}$  as well as  $Sh_x Re_x^{-1/2}$ . However, nanoparticle Cu dominated the results obtained in Figure 6. The value of  $f''(0)$  for

**Table 2.** Value of  $f''(0)$  for several values of  $\epsilon$  when  $\phi = 0.1$

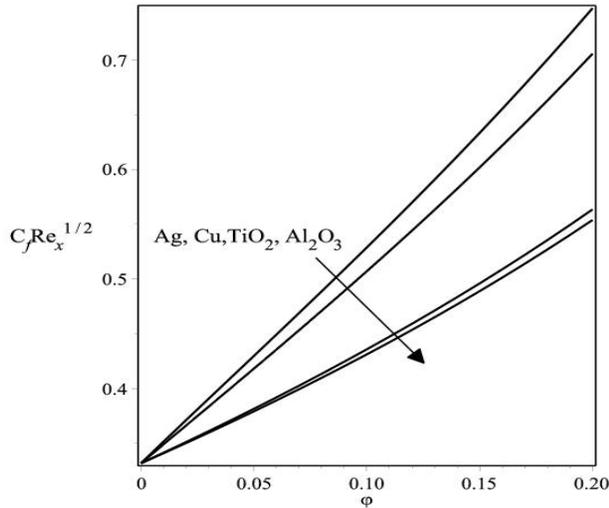
$\epsilon$	$f''(0)$			
	Ag	Cu	$TiO_2$	$Al_2O_3$
0	1.510003	1.447977	1.244317	1.231074
0.5	0.873834	0.837940	0.720083	0.712419
1	0	0	0	0
2	-2.312079	-2.217105	-1.905267	-1.884989



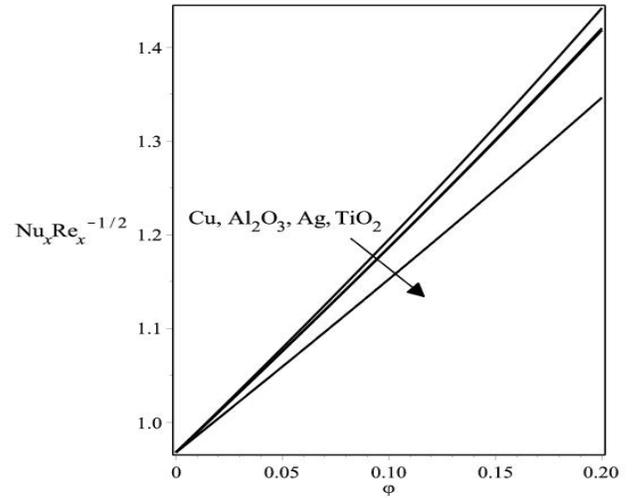
**Figure 3.** Variation of  $Nu_x Re_x^{-1/2}$  with  $Nt$  for various nanoparticles when  $\epsilon = 1.2$ ,  $Pr = 6.2$ ,  $\phi = 0.1$ ,  $Le = 3$  and  $Nb = 0.1$ .



**Figure 4.** Variation of  $Sh_x Re_x^{-1/2}$  with  $Nt$  for various nanoparticles when  $\epsilon = 1.2$ ,  $Pr = 6.2$ ,  $\phi = 0.1$ ,  $Le = 3$  and  $Nb = 0.1$ .



**Figure 5.** Variation of  $C_f Re_x^{1/2}$  with  $\phi$  for various nanoparticles when  $Nb = Nt = 0.1$ ,  $\epsilon = 1.2$ ,  $Pr = 6.2$  and  $Le = 3$ .

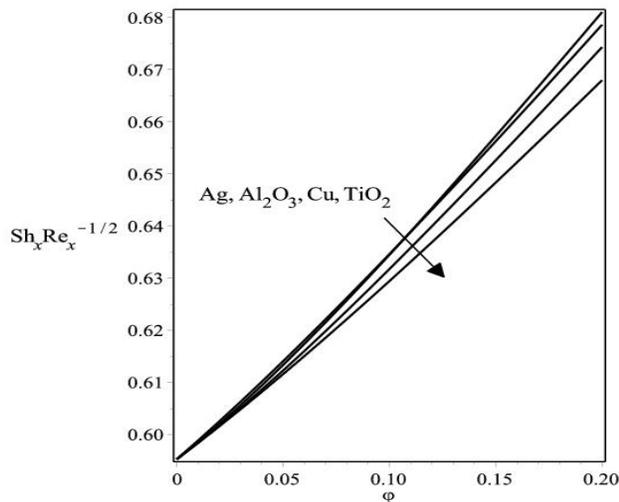


**Figure 6.** Variation of  $Nu_x Re_x^{-1/2}$  with  $\phi$  for various nanoparticles when  $Nb = Nt = 0.1$ ,  $\epsilon = 1.2$ ,  $Pr = 6.2$  and  $Le = 3$ .

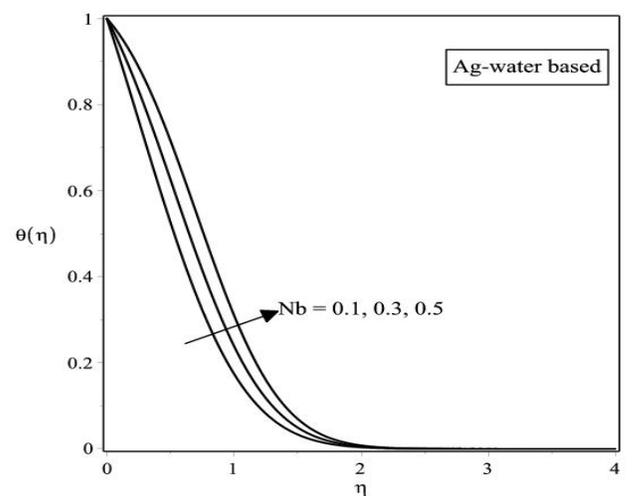
different  $\epsilon$  for Ag, Cu,  $TiO_2$  and  $Al_2O_3$  in water when  $Pr = 6.2$ ,  $Nt = 0.1$ ,  $Le = 3$ ,  $Nb = 0.1$  and  $\phi = 0.1$  is shown in Table 2.

Figure 8 illustrate the changes of  $Nb$  on the temperature profile  $\theta(\eta)$ , mean while Figure 9 shows the changes of nanoparticle concentration profile  $\phi(\eta)$  for Ag-water fluid when  $\epsilon = 1.2$ ,  $Pr = 6.2$ ,  $\phi = 0.1$ ,  $Le = 3$  and  $Nt = 0.1$ .

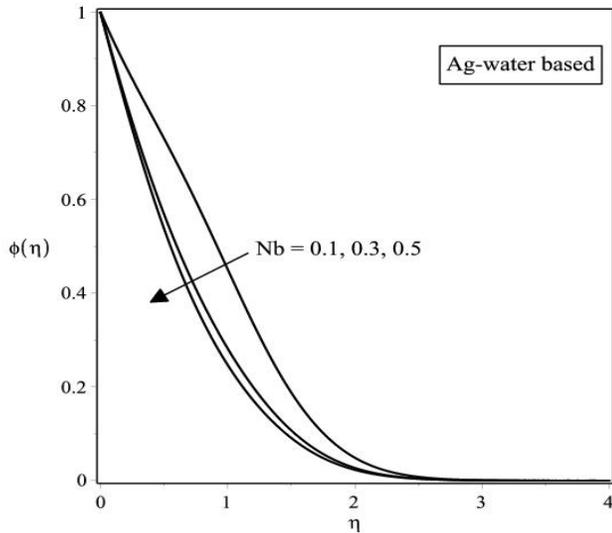
It is noticed that a gradual increase of the temperature profile and but decreases the nanoparticle concentration profile with increasing  $Nb$ . Hence, in Figure 8, the boundary layer thickness increases when  $Nb$  increases. The opposite trend is found in Figure 9, where the boundary layer thickness reduces when  $Nb$  rise up. The increasing of Brownian motion is able to raise the thermal state of



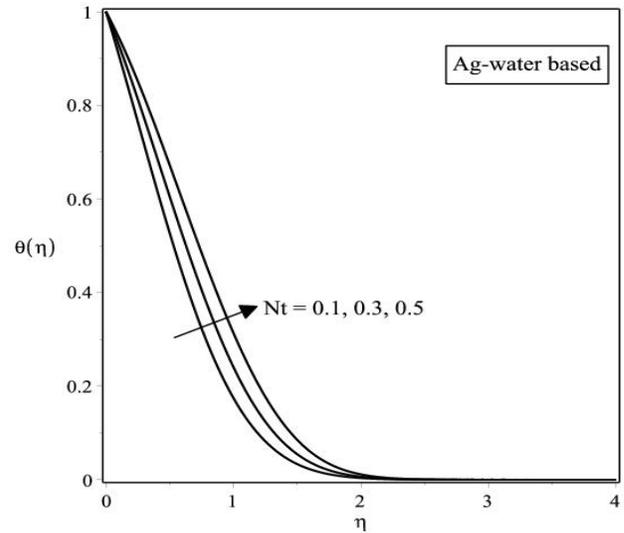
**Figure 7.** Variation of  $Sh_x Re_x^{-1/2}$  with  $\phi$  for various nanoparticles when  $Nb = Nt = 0.1$ ,  $\epsilon = 1.2$ ,  $Pr = 6.2$  and  $Le = 3$ .



**Figure 8.** Temperature profile for various  $Nb$  when  $\epsilon = 1.2$ ,  $Pr = 6.2$ ,  $\phi = 0.1$ ,  $Le = 3$  and  $Nt = 0.1$ .



**Figure 9.** Concentration profile for various  $Nb$  when  $\epsilon = 1.2$ ,  $Pr = 6.2$ ,  $\phi = 0.1$ ,  $Le = 3$  and  $Nt = 0.1$ .



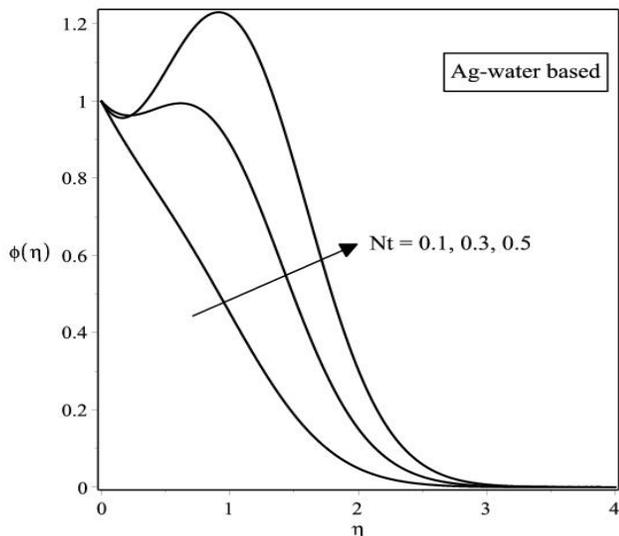
**Figure 10.** Temperature profile for various  $Nt$  when  $\epsilon = 1.2$ ,  $Pr = 6.2$ ,  $\phi = 0.1$ ,  $Le = 3$  and  $Nb = 0.1$ .

the nanofluid due to collision of the nanoparticles and based fluid.

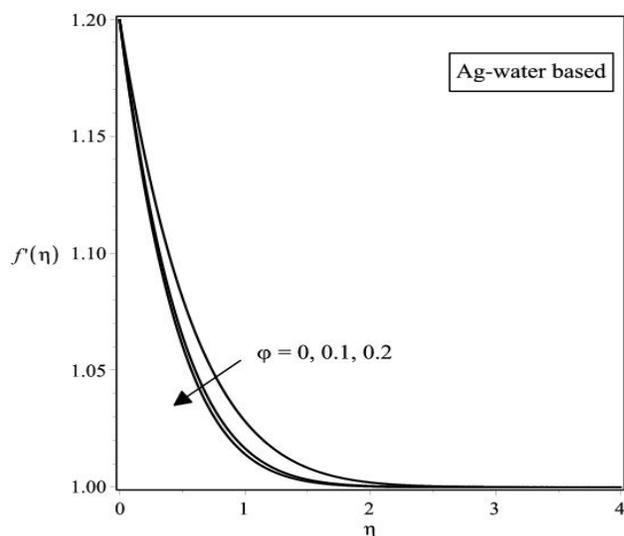
The effects of  $Nt$  on  $\theta(\eta)$  profile and  $\phi(\eta)$  profile for Ag-water fluid when  $\epsilon = 1.2$ ,  $Pr = 6.2$ ,  $\phi = 0.1$ ,  $Le = 3$  and  $Nb = 0.1$  are depicted in Figures 10 and 11, respectively. The figures demonstrate that  $\theta(\eta)$  and  $\phi(\eta)$  profiles

increase slightly when  $Nt$  increases. The increasing  $Nt$ , the boundary layer thicknesses are more thicker. The smaller  $Nt$ , the thermophoresis effect blows the boundary layer thicknesses apart from the surface.

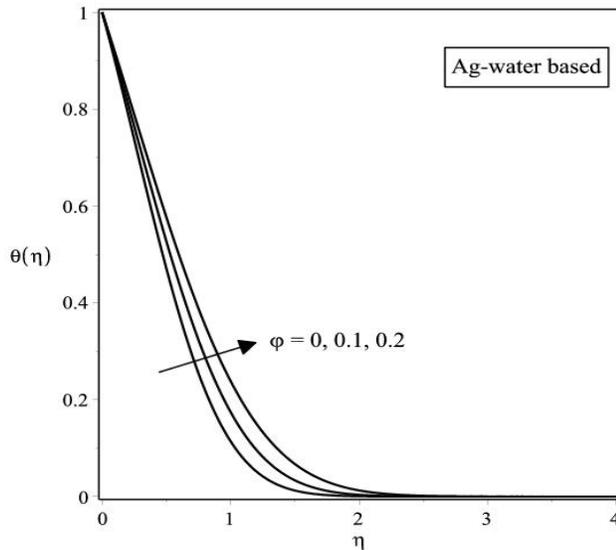
Figures 12, 13 and 14 demonstrate the effects of  $\phi$  on  $f'(\eta)$ ,  $\theta(\eta)$  and  $\phi(\eta)$  profiles for Ag-water fluid when  $\epsilon =$



**Figure 11.** Concentration profile for various  $Nt$  when  $\epsilon = 1.2$ ,  $Pr = 6.2$ ,  $\phi = 0.1$ ,  $Le = 3$  and  $Nb = 0.1$ .



**Figure 12.** Velocity profile for various  $\phi$  when  $\epsilon = 1.2$ ,  $Pr = 6.2$ ,  $Le = 3$  and  $Nb = Nt = 0.1$ .

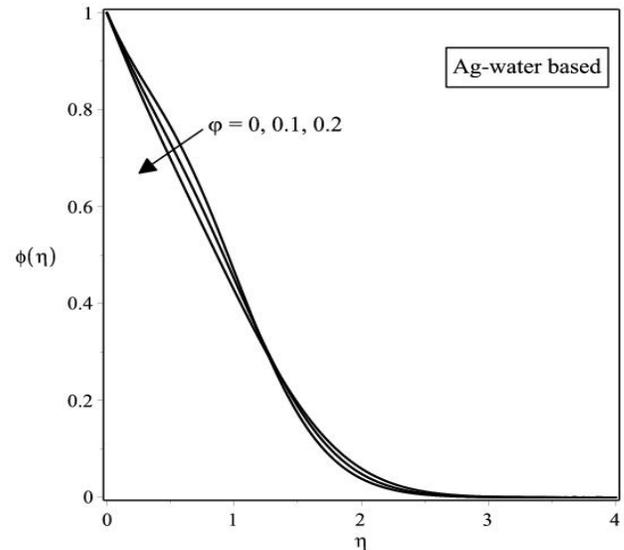


**Figure 13.** Temperature profile for various  $\phi$  when  $\varepsilon = 1.2$ ,  $Nt = 0.1$ ,  $Pr = 6.2$ ,  $Le = 3$  and  $Nb = 0.1$ .

1.2,  $Nb = Nt = 0.1$ ,  $Le = 3$  and  $Pr = 6.2$ . From Figure 12 and Figure 14, when  $\phi$  increases,  $f'(\eta)$  and  $\phi(\eta)$  profiles reduce. Meanwhile in Figure 13, the reverse behaviour are found on  $\theta(\eta)$  profile. The larger nanoparticle volume fraction causes the nanofluid more viscous and the mixture of convection is weaker, thus it will reduce the local Nusselt number due to high viscosity of nanofluid. The numerical results are supported because for Figure 17 till Figure 23, we found that all these profiles satisfy the boundary conditions Equation (11) asymptotically.

## 4. Conclusions

The nanofluid flow past a stretching sheet was examined. The influences of parameters  $Nb$ ,  $Nt$ ,  $\phi$  and  $\varepsilon$  on, and were studied. From the investigation, a unique solution is seen for stretching sheet. Otherwise, when the  $\phi$  increases,  $Nb$  decreases and  $Nt$  decreases, the heat transfer rate elevated but reduces the mass transfer rate. Besides, the heat transfer rate reduces meanwhile the mass transfer rate elevated when  $\varepsilon$  increases.



**Figure 14.** Concentration profile for various  $\phi$  when  $\varepsilon = 1.2$ ,  $Nt = 0.1$ ,  $Pr = 6.2$ ,  $Nb = 0.1$  and  $Le = 3$ .

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